Interaction Mechanism Insights on the Solvation of Fullerene B₈₀ with Choline-based Ionic Liquids

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Supporting Information

ABSTRACT: Beyond carbon allotropes, other nanostructures such as fullerene B₈₀ are attracting a growing interest due to their potential applications. The use of new materials based on fullerene B₈₀ is still in a premature stage; however, many of these applications would require the use of B₈₀ in solution. This paper reports an unprecedented density functional theory (DFT) analysis on the interaction mechanism between B₈₀ and two choline-based ionic liquids as a first insight for the fullerene B₈₀ solvation by ionic liquids. The analysis of properties such as binding energies, charge distributions or intermolecular interactions shed light on the main features, which should govern interaction between ionic liquids and fullerene B₈₀. In addition, the optimization of systems composed by six ionic pairs around a fullerene B₈₀ has supplied some information about the first solvation shell at the molecular level. As a summary, this paper provides the first insights in the rational design of ionic liquids with suitable properties for the solvation of B₈₀.

1. INTRODUCTION

There has been a growing interest in carbon nanostructures, such as fullerenes, nanotubes, and graphene flakes or sheets due to their promising potential as building blocks in a wide range of applications, such as biotechnology, nanobiomedicine, energy storage, and electronic devices. Because of their significant promising structural, electronic, and chemical properties, there is also a rising attention in the search for alternative nanostructures that are more suitable for further applications on advanced cutting edge research fields as mentioned above. Boron, which has only one p electron, is a neighboring element with carbon in the periodic table, thus it has been prompted that boron could also form nanostructures (or, in general, boron clusters). Boron atom also tends to form sp² hybridization, while its electron dispersion (or, in general, boron clusters). Boron atom also tends to form sp² hybridization, while its electron dispersion.

Among these new boron nanostructures, hollow fullerene B₈₀ has attracted considerable interest due to the existence of a stable cage B₈₀ with icosahedral symmetry (T₅₂). Later, other alternative B₈₀ geometries with tetrahedral symmetry (T₄₂) and core−shell structures have been proposed. This T₄₂ symmetry seems to be thermodynamically more stable. However, Sadrzadeh et al. found that T₄₂ symmetry is obtained by slight distortion of T₅₂ structure. In fact, T₅₂ and T₄₂ structures (and other one with C₁ symmetry) are close in energy and have almost identical structure. Although B₈₀ has not yet been manufactured, theoretical simulations point out to a large cohesive energy (∼5.76 eV/atom), high deformation temperature and high formalization stability. Meanwhile, a large number of computational works on its electronic structure, stability and reactivity have been studied along with various applications, including hydrogen storage or CO₂ capture, have been published. Despite all these theoretical simulations, fullerene B₈₀ still has unexplored applications in a liquid media. For instance, the successful application of its analogue fullerene C₆₀ requires a deep understanding on its behavior in the solution state, i.e., C₆₀ solvation. To our knowledge, studies on the solvation of fullerene B₈₀ has not been reported in the open literature up to date.

Among the plethora of solvents, ionic liquids (IL) have proven to be suitable alternatives to traditional solvents due to their unique features such as good thermal and chemical stability, non flammability and almost null vapor pressure. Previous studies were pointed out that the solvation shell around the C₆₀ fullerene is mainly characterized...
by ion-C_{60} π−π stacking interactions. Yet, further analysis on the interaction mechanism between 24 IL and C_{60} through density functional theory simulations revealed that C_{60} solvation could be also achieved by ILs with deep HOMO energy level and weak interionic interactions as well. In this sense, Jang et al. studied C_{60} solvation by water molecules, which revealed that negative fullerene surface potential along hydrogen bonding network around C_{60} is the main driving force for C_{60} self-aggregation in water.

At the molecular level, the solvation capability of ILs would be related with the strength of IL–B_{80} interactions, i.e. strengthening interaction energy is needed for stabilizing B_{80} fullerene in the solvent. Thus, interaction energy between B_{80} and ILs, estimated thorough DFT simulations, can be used as an approximation on inferring the main IL features at the molecular level needed for the design of task-specific ILs for B_{80} solvation. Hence, a DFT study on the interaction mechanism between selected ILs and fullerene B_{80} is reported in this work. Choline lactate ([CH][LAC]) and choline benzoate ([CH][BE]) ILs were selected in this work (Figure 1). Choline cation-based ILs are a new generation of green solvents ([CH][LAC]) and choline benzoate ([CH][BE]) ILs are a new generation of green solvents.41,42 Likewise, the combination of [CH] cation with favorable properties such as null toxicity and high biodegradability.43 ([CH][LAC]) and [CH][BE] ILs were previously studied by our group through DFT simulations and their stable configurations were optimized at PBE-D2/DZP level. Once the B_{80} was fully optimized, one ionic pair was introduced to the simulations. In order to determine the most favorable position of the ionic pair, different structures were assessed as starting geometries (Figure S1, Supporting Information) for energy minimization. On the basis of the most favorable IL disposition onto the B_{80} surface, IL_{n}–B_{80} clusters composed by n ionic pairs (n = 2, 3, 4, 6) around one B_{80} were also optimized.

Atoms in molecules (AIM) theory and the analysis of the reduced density gradient (RDG) at low densities have been used for better understanding of the nature of the interaction mechanism. According to AIM theory, the existence of an interaction is indicated by the presence of a bond critical point (BCP), which can be featured based on its electronic density (ρ) and laplacian values (V^2ρ). Similarly, ring/cage structures were characterized by their corresponding ring/cage critical point (RCP/CCP). Aimed at clarifying data analysis, attention has been mainly given on the bond critical points, while some information regarding to the adsorption process could also be obtained from ring and cage critical points. RDG analysis is based on electron density descriptor to identify the main interactions, which is based on the visualization of RDG isosurfaces. The isosurfaces of the reduced density gradient are defined as s = (1/(2(3π)^1/3))(((Vρ)/(|V^2ρ|)^1/3)); Strength and nature of the interactions is determined through the sign of the second density Hessian eigenvalue. Thus, by means of a color mapping scheme (based on the value of the electronic density times of the second eigenvalue of the Hessian), strong nonbonded overlaps appear as red regions, strong intermolecular interactions are plotted as a blue localized regions, while weak contacts (dispersive forces) appear as large green regions. Charge transfer process were determined using the ChelpG method, which has proven to be adequate for describing charge distributions in ILs. AIM and RDG analysis were carried out using MultiWFN code, which was also used to compute ChelpG atomic charges.

2. THEORETICAL DETAILS

DFT simulations were carried out using PBE functional as implemented in the SIESTA 3.2 package, along norm-conserving Troullier–Martins pseudopotentials and numerical double-ζ polarized (DZP) basis sets. The PBE functional has been proven to perform well for large systems with a relatively low computational cost. Ionic pair and IL–B_{80} interactions would involve a large component of noncovalent dispersion interactions. It is a well-known situation that dispersion interactions are not satisfactory described by ordinary DFT methods, and thus, long-range corrections using the Grimme’s scheme were added to PBE functional (PBE-D2) and this method has been successfully applied to study pristine fullerene B_{80} and systems related to B_{80}. All calculations were done with an energy mesh cutoff of 400 Ry, while structural relaxations were done by conjugate gradients, with convergence criteria enforced on all atoms which do not exceed 0.04 eV/Å.

Initial structure for B_{80} structure was obtained from the work of Muya et al. and its optimization studies built on this extracted model. The optimized molecular structure in this work were compared with the structure reported by Muya et al. and is consistent with this previous study. [CH][LAC] and [CH][BE] ILs were previously studied by our group through DFT simulations and their stable configurations were optimized at PBE-D2/DZP level. Once the B_{80} was fully optimized, one ionic pair was introduced to the simulations. In order to determine the most favorable position of the ionic pair, different structures were assessed as starting geometries (Figure S1, Supporting Information) for energy minimization. On the basis of the most favorable IL disposition onto the B_{80} surface, IL_{n}–B_{80} clusters composed by n ionic pairs (n = 2, 3, 4, 6) around one B_{80} were also optimized.

3. RESULTS AND DISCUSSION

The study of IL–B_{80} systems is divided into three sections aimed at assessing the main features of the interaction between fullerene B_{80} and choline-based ILs. The first section analyzes the main properties related with the interaction between ions, which are need for a deeper understanding on the interaction mechanism between IL and B_{80}. The second section provides a
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Table 1. Main Molecular Parameters of the Most Stable Optimized Structures of [CH][LAC] and [CH][BE] Ionic Liquids\(^a\)

<table>
<thead>
<tr>
<th></th>
<th>[CH][LAC]</th>
<th></th>
<th></th>
<th>[CH][BE]</th>
<th></th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>length/Å</td>
<td>ρ/au</td>
<td>(V_\rho^2/au)</td>
<td>length/Å</td>
<td>ρ/au</td>
<td>(V_\rho^2/au)</td>
</tr>
<tr>
<td>intermolecular interactions</td>
<td></td>
<td></td>
<td></td>
<td></td>
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</tr>
<tr>
<td>(a_{10}) (O–H)</td>
<td>1.863</td>
<td>0.0360</td>
<td>0.1228</td>
<td>(a_{23}) (O–H)</td>
<td>1.673</td>
<td>0.0668</td>
</tr>
<tr>
<td>(a_{11}) (O–H)</td>
<td>1.512</td>
<td>0.0777</td>
<td>0.1122</td>
<td>(a_{22}) (O–H)</td>
<td>1.980</td>
<td>0.0265</td>
</tr>
<tr>
<td>(a_{12}) (O–H)</td>
<td>2.471</td>
<td>0.0111</td>
<td>0.0400</td>
<td>(a_{23}) (O–H)</td>
<td>2.122</td>
<td>0.0213</td>
</tr>
<tr>
<td>(a_{13}) (O–H)</td>
<td>2.029</td>
<td>0.0247</td>
<td>0.0777</td>
<td>(a_{24}) (O–H)</td>
<td>1.816</td>
<td>0.0382</td>
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<td>(a_{14}) (O–H)</td>
<td>1.883</td>
<td>0.0309</td>
<td>0.1095</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>dihedral angles/degrees</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
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</tr>
<tr>
<td>(τ_{10}) (CCOH)</td>
<td>~1.0</td>
<td></td>
<td></td>
<td>(τ_{23}) (CCOH)</td>
<td>132.3</td>
<td></td>
</tr>
<tr>
<td>(τ_{11}) (CCOH)</td>
<td>137.3</td>
<td></td>
<td></td>
<td>(τ_{21}) (OCCH)</td>
<td>0.0</td>
<td></td>
</tr>
<tr>
<td>(τ_{12}) (OCCC)</td>
<td>1.1</td>
<td></td>
<td></td>
<td></td>
<td></td>
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</tr>
<tr>
<td>(q^{+}/e) &amp; (q^{-}/e)</td>
<td>0.76 &amp; 0.69</td>
<td></td>
<td></td>
<td>0.69 &amp; 0.00</td>
<td></td>
<td></td>
</tr>
<tr>
<td>(\text{BE}_{\text{IL}}/\text{kcal mol}^{-1})</td>
<td>87.18 (4.59)</td>
<td></td>
<td></td>
<td>90.20 (4.60)</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\(^a\)AIM parameters related with intermolecular interactions (electronic density, \(ρ\), and its laplacian, \(V_\rho^2\)) as well as total charge over choline (\(q^{+}\)) and lactate or benzoate (\(q^{-}\)) ions computed according ChelpG scheme are also collected. See Figure 1 for labeling. \(^b\)For isolated ionic liquids \(q^{+} = -q^{-}\). Values in parentheses stand for dispersion contribution to the total binding.

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Detailed analysis of the interaction mechanism of IL–B\(_{80}\) systems. Therefore, properties such as binding energies, intermolecular interactions and electronic structure have been studied. The last section reports and insight into the first solvation shell of B\(_{80}\) by ILs. For this, DFT simulations of (IL)\(_{n}\)–B\(_{80}\) \((n = 2, 3, 4, 6)\) have been also carried out.

3.1. [CH][LAC] and [CH][BE] Ionic Pairs. The main features of [CH][LAC] and [CH][BE] ILs have been previously studied by our group,\(^{57,58}\) therefore in this work only the main features of these materials were analyzed and discussed. Figure 1 (right) shows optimized structures of the most stable conformations of both [CH][LAC] and [CH][BE] ILs at PBE-D2/DZP theoretical level, while Table 1 gathers the main parameter related with interactions between ions. Interionic binding energies (\(\text{BE}_{\text{IL}}\)) were used as a measurement of the interaction strength between ions and it is defined as

\[
\text{BE}_{\text{IL}} = (E_{\text{cat}} + E_{\text{an}}) - E_{\text{IL}},
\]

where \(E_{\text{cat}}\), \(E_{\text{an}}\), and \(E_{\text{IL}}\) stand for the total energy of the cation, anion, and ionic pair, respectively. Estimated \(\text{BE}_{\text{IL}}\) values are reported as 87.18 kcal mol\(^{-1}\) for [CH][LAC] and 90.20 kcal mol\(^{-1}\) for [CH][BE] (Table 1). Charge distributions were calculated according to the ChelpG model, and those results showed that the charge of ions is 0.76\(e^−/0.69\text{e}^-\), in absolute value, for [CH][LAC]/[CH][BE] ILs, respectively, and thus showing an important charge transfer from the anion to the [CH] cation equivalent to 0.24 \(\text{e}^-/0.31 \text{e}^−\), which is in agreement with the highest \(\text{BE}_{\text{IL}}\) value of [CH][BE] ionic liquid. Dispersion contributions (such as hydrogen bonds) also represent a small contribution to the total binding energy between ions. The optimized geometries show the presence of intermolecular hydrogen bonds between [CH]\(^+\) cation and [LAC]\(^−\) or [BE]\(^−\) anions, through the cation hydroxyl and carboxylate groups in the corresponding ions (labeled as \(a_{12}\) or \(a_{21}\), respectively). Likewise, an additional weak intermolecular hydrogen bond between COO\(^−\) group and H atoms in [CH] cation are also found (\(a_{13}−a_{14}\) and \(a_{23}−a_{24}\) for [CH][LAC] and [CH][BE] respectively). On the basis of shorter intermolecular bond length and AIM features (Table 1), intermolecular hydrogen bond between [LAC]\(^−\) and [CH]\(^+\) (\(a_{13} = 1.512, \rho = 0.0777 \text{ au}\)) is slightly stronger than for [BE]\(^−\) and [CH]\(^+\) (\(a_{23} = 1.673, \rho = 0.0668 \text{ au}\)). Nonetheless, the sum of the electronic density overall intermolecular interactions yields values slightly larger (0.0084 au) for [CH][BE]. Table 1 also shows the dispersion energy contribution (\(\text{BE}_{\text{dis},\text{IL}}\)) to the total binding energy values of around 4.60 kcal mol\(^{-1}\) for both ILs, which is quantified according to Grimme’s approach for the PBE-D2 functional.\(^{52}\)

3.2. IL–B\(_{80}\) Systems: Key Features of the Interaction Mechanism. Aimed at obtaining the most favorable disposition between selected ILs and fullerene B\(_{80}\) different geometries were used as starting point for the optimizations (see Figure S1 and theoretical details in the Supporting Information). Among all essayed structures, those of minimal energy were further used for the analysis of IL–B\(_{80}\) systems. The most favorable configurations of ionic pairs interacting with B\(_{80}\) are reported in Figures 2 and 3, while the main parameters (intermolecular bond lengths and their AIM features, charge distributions, and binding energies) are collected in Table 2. The main features of ILs adsorbed on the surface of fullerene B\(_{80}\) is discussed herein. Several authors have pointed out that the change on cation–anion interaction strength can be used to understand the entropic contribution to

Figure 2. Side (up) and top (bottom) views for optimized structure at PBE-D2/DZP level corresponding to [CH][LAC]–B\(_{80}\) system along structural parameters related with intermolecular interactions (left) and RDG isosurfaces (right), whose green color indicates van der Waals interactions. Red and green points stand for RCP and CCP, respectively, related with intermolecular interactions. BCP were omitted for simplicity.
Since coulomic forces stand for the main forces between ions in the ionic liquid bulk. BE\textsubscript{\textsc{il}} values have been calculated for both ILs but using their geometries taken from IL–B\textsubscript{80} systems (Table 2). Both ILs diminish their binding energies up to 76.59 kcal mol\textsuperscript{-1} (energy diminution is equal to 10.59 kcal mol\textsuperscript{-1})/76.70 kcal mol\textsuperscript{-1} (13.50 kcal mol\textsuperscript{-1}) for [\text{CH}][\text{LAC}]/[\text{CH}][\text{BE}]. For [\text{CH}][\text{LAC}] IL, the weakening in the interaction between ions is mainly due to two factors: (i) the small charge transfer from the surface anion to the surface of B\textsubscript{80} (0.07 e\textsuperscript{−}) hinders the charge transfer process between ions, which is equal to 0.15 e\textsuperscript{−} (anion–cation charge transfer is equal to 0.24 e\textsuperscript{−} for the isolated ionic pair); (ii) softening of the intermolecular hydrogen bonds between ions (which can be assessed through the elongation of the intermolecular lengths as well as lower electronic density values), where even \(\varepsilon\textsubscript{11}\) was not found. The main interaction between both ions (\(\varepsilon\textsubscript{11}\)) is elongated (0.161 Å) as well as its electronic density decreases, in agreement with a lower dispersion contribution (3.26 kcal mol\textsuperscript{-1}) in comparison to 4.59 kcal mol\textsuperscript{-1} for isolated [\text{CH}][\text{LAC}]. Regarding the [\text{CH}][\text{BE}–B\textsubscript{80} system, the [\text{BE}] anion becomes slightly less negative, which allows a charge transfer to the B\textsubscript{80} surface (0.07 e\textsuperscript{−}) while charge transfer (CT) between ions is not importantly affected (0.30 e\textsuperscript{−}). The presence of fullerene B\textsubscript{80} carries the break of \(\varepsilon\textsubscript{12}\) and \(\varepsilon\textsubscript{23}\) intermolecular bond, while \(\varepsilon\textsubscript{21}\) and \(\varepsilon\textsubscript{32}\) slightly decrease (despite their bond lengths are shortened). As seen below, the largest charges on BE\textsubscript{\textsc{il}} values for [\text{CH}][\text{BE}] ionic liquids could be related to stronger IL–B\textsubscript{80} interactions.

Table 2. Main Molecular Parameters for the Most Stable Structures Optimized at PBE-D2/DZP Levels of IL–B\textsubscript{80} Systems

<table>
<thead>
<tr>
<th>IL</th>
<th>(\varepsilon\textsubscript{\textsc{il}}) length/Å</th>
<th>(\rho)/e/Å (\nabla^2\rho)/e/Å</th>
<th>(\varepsilon\textsubscript{\textsc{il}}) BE\textsubscript{\textsc{il}}/kcal mol\textsuperscript{-1}</th>
</tr>
</thead>
<tbody>
<tr>
<td>[\text{CH}][\text{LAC}]</td>
<td>1.995</td>
<td>0.0273</td>
<td>0.1070</td>
</tr>
<tr>
<td>[\text{CH}][\text{BE}]</td>
<td>1.668</td>
<td>0.0500</td>
<td>0.1387</td>
</tr>
</tbody>
</table>

\(\varepsilon\textsuperscript{11}\) is the sum of \(\varepsilon\textsuperscript{\textsc{il}}\) for isolated ions. \(\varepsilon\textsuperscript{\textsc{il}}\) BE\textsubscript{\textsc{il}}/kcal mol\textsuperscript{-1} represents the binding energy, \(\varepsilon\textsuperscript{12}\) and \(\varepsilon\textsuperscript{23}\) are the sum of \(\varepsilon\textsuperscript{\textsc{il}}\) for ions at the intermolecular interaction.

Since coulomic forces stand for the main forces between ions in the ionic liquid bulk. BE\textsubscript{\textsc{il}} values have been calculated for both ILs but using their geometries taken from IL–B\textsubscript{80} systems (Table 2). Both ILs diminish their binding energies up to 76.59 kcal mol\textsuperscript{-1} (energy diminution is equal to 10.59 kcal mol\textsuperscript{-1})/76.70 kcal mol\textsuperscript{-1} (13.50 kcal mol\textsuperscript{-1}) for [\text{CH}][\text{LAC}]/[\text{CH}][\text{BE}]. For [\text{CH}][\text{LAC}] IL, the weakening in the interaction between ions is mainly due to two factors: (i) the small charge transfer from the surface anion to the surface of B\textsubscript{80} (0.07 e\textsuperscript{−}) hinders the charge transfer process between ions, which is equal to 0.15 e\textsuperscript{−} (anion–cation charge transfer is equal to 0.24 e\textsuperscript{−} for the isolated ionic pair); (ii) softening of the intermolecular hydrogen bonds between ions (which can be assessed through the elongation of the intermolecular lengths as well as lower electronic density values), where even \(\varepsilon\textsubscript{11}\) was not found. The main interaction between both ions (\(\varepsilon\textsubscript{11}\)) is elongated (0.161 Å) as well as its electronic density decreases, in agreement with a lower dispersion contribution (3.26 kcal mol\textsuperscript{-1}) in comparison to 4.59 kcal mol\textsuperscript{-1} for isolated [\text{CH}][\text{LAC}]. Regarding the [\text{CH}][\text{BE}–B\textsubscript{80} system, the [\text{BE}] anion becomes slightly less negative, which allows a charge transfer to the B\textsubscript{80} surface (0.07 e\textsuperscript{−}) while charge transfer (CT) between ions is not importantly affected (0.30 e\textsuperscript{−}). The presence of fullerene B\textsubscript{80} carries the break of \(\varepsilon\textsubscript{12}\) and \(\varepsilon\textsubscript{23}\) intermolecular bond, while \(\varepsilon\textsubscript{21}\) and \(\varepsilon\textsubscript{32}\) slightly decrease (despite their bond lengths are shortened). As seen below, the largest charges on BE\textsubscript{\textsc{il}} values for [\text{CH}][\text{BE}] ionic liquids could be related to stronger IL–B\textsubscript{80} interactions.
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The strength of IL–B80 interactions is related to the suitability of ILs as B80 solvents. Binding energies of IL–B80 systems (BE_{IL–B80}) were defined as

$$BE_{IL–B80} = (E_{B80} + E_{IL}) - E_{IL–B80}$$  

where $E_{B80}$, $E_{IL}$, and $E_{IL–B80}$ are the total energy of the fullere, ionic pair, and IL–B80, respectively. As seen in Table 2, binding energies lie between 43.86 kcal mol$^{-1}$ for [CH][LAC] and 45.95 kcal mol$^{-1}$ for [CH][BE], where dispersion contributions (36.98% and 47.97% for [CH][LAC] and [CH][BE] respectively) would be one of the main factors.

Interactions between [CH]-based ILs and fullere B80 have been described through the topological analysis of the electronic density (AIM theory) and RDG isosurfaces, Figures 2 and 3. Green color of the regions between the IL and B80 molecule points out van der Waals interactions as the main driving force. Aimed at studying all BCPs as a whole, we have defined the sum of the electronic density ($\sum_{BCP}$) over all BCPs related with cation/anion–B80 interactions: $\sum_{BCP} = \sum_{BCP,cat} + \sum_{BCP,ani}$ (see Table 2). Instead of the selected ILs, cation–B80 interactions take place through hydrogen bonds between H atoms of methyl groups and negatively charged boron atoms, with bond lengths equal to (in average) 2.491 Å/2.462 Å for [CH][LAC]/[CH][BE] ILs, while $\sum_{BCP,cat}$=0.0334 au/0.0484 au.

The main feature of anion–B80 interactions is characterized by the presence of a chemical bond between COO$^{-}$ group and B80 surface ($c_{11}/c_{12}$ for [CH][LAC]/[CH][BE]). The laplacian of the electronic density ($V^2\rho$) can be used to classify different bonds as $V^2\rho > 0$ the bond is defined as closed-shell interaction (which could be an ionic, hydrogen or van der Waals bond); (ii) $V^2\rho < 0$ the bond is defined as shared interaction (i.e., a covalent bond). For both ionic liquids, O–B interaction between carboxylate moiety and B80 surface yields a $V^2\rho$ values larger than zero. In addition, the absence of RDG isosurface between both O and B atoms as well as the (small) charge transfer from the anion to the B80 surface would point out to the presence of an ionic bond between the COO$^{-}$ moiety and the B80 surface. In addition, both bonds yield electronic density values of around 0.1363 au, which is in agreement with values obtained for other systems with typical ionic bonds.

Although both $c_{11}$ and $c_{12}$ bonds show similar features, there are differences in anion–B80 interactions. In addition to the O–B bond ($c_{11}$), [LA]–B80 interactions are set up between H atoms located in methyl groups ($c_{12}$ and $c_{13}$) and the B80 surface, with a bond length of around 2.498 Å. Regarding to [BE]–B80 interactions, there is a clear $\pi$-stacking between phenyl moiety and the B80 surface, with a distance of around 3.0 Å. Palusiak et al. found that ring critical point features (mainly the electronic density, $\rho$) are connected with the $\pi$-electronic delocalization. Then, we have conjectured a relationship of electronic delocalization between IL and B80 and B80. Although both ILs yield similar $\sum_{BCP,ani}$ values, larger values of $\sum_{BCP}$ for [CH][BE] IL agrees with a higher electronic delocalization due to the $\pi$-stacking as well as with a larger dispersion contribution to the total BE_{IL–B80}. Therefore, the greatest BE_{IL} diminution for [CH][BE] is related to stronger IL–B80 interactions.

For simplicity, discussions on the B80 structure were not deeply done herein as they have been extensively studied in the literature by itself.31,32,33,34,35,36 The presence of a chemical bond between COO$^{-}$ group and B80 leads to a small distorsion of B80 symmetry, which has not dramatic effects on the features of fullere B80. For instance, as seen bellow, density of states of pristine B80 is very similar than partial density of states from B80 in IL–B80 systems. Another important feature is the atomic charge distribution of B80. Published works assessing charge distribution of B80 are based on Mulliken atomic charges.24,25,26,27 ChelpG charge distribution of pristine B80 leads to similar qualitative results, i.e. boron atoms in the center of hexagonal faces own positive charge, whereas frame boron atoms are negatively charged. The bond between COO$^{-}$ group and boron surface is carried out with a positive charged boron atom located in the center of a hexagonal face (see Figures 2 and 3). Nonetheless, hydrogen bond labeled as $c_{31}$ in [CH][LAC]/[CH][BE] system brings that a boron atom located in the center of a hexagonal face becomes negative charged. Although this is a very weak bond (see AIM features in Table 2), this new charge distribution could promote some B80 stability problems.

The results reported in this section confirm that [CH][LAC] and [CH][BE] ionic liquids would distort B80 geometry, which could be a problem for the use of B80 dispersed in ILs. The selection of these ILs were based on considering the absence of previous studies and targetted to minimize both the economic and environmental issues that are related with the use of classical ILs. Nevertheless, the results reported in this work reveal useful conclusions for the search of task-specific ILs regarding to B80 solvation. Ions with aromatic motifs could be adequate as B80 solvents due to strong interactions through $\pi$–$\pi$ stacking; whereas anions with functional groups, which tend to allocate much of its negative charge (e.g., COO$^{-}$), shall be avoided.

The main features of the electronic structure of IL–B80 systems were also studied. Figure S2 (Supporting Information) shows the total density of states (DOS) of pristine B80 and IL–B80 systems as well as the partial density of states (PDOS) corresponding to B80 contributions for IL–B80 systems. The total density of states of pristine B80 and the partial density of states corresponding to B80 atoms in IL–B80 systems own similar contours. Therefore, interactions with the selected ILs do not have large effects on the electronic structure of fullerene B80. Figure 4 shows the partial density of states from both ions for isolated ILs and both IL–B80 systems as well as the molecular orbital contours for the highest occupied and lowest unoccupied molecule orbitals (HOMO and LUMO) of the ILs. For isolated [CH] and[LA]ILs, HOMO/LUMO levels are mainly delocalized over [LA] anion/[CH] cation, with a HOMO–LUMO energy gap $\Delta E_{H-G} = 3.96$ eV. Regarding [CH][BE] IL, the HOMO orbital is mainly over carboxylate group with some contribution from the cation, while the LUMO level is mainly localized over the [BE] anion, where $\Delta E_{H-G} = 3.63$ eV. Because of the interaction with B80, LUMO orbital becomes LUMO+1, whose energy destabilization is equal to 0.37 eV/1.10 eV for [LA]/[BE] anions. Both HOMO and LUMO orbitals corresponding to the ionic liquid are delocalized over the anion, although they also own an important contribution from B80 fullere, which agrees with the charge transfer process between the ionic liquid and B80 above-described. Both HOMO energies are of around $-6.37$ eV, while $\Delta E_{H-G} = 4.28$ eV/3.70 eV for [CH][LAC]/[CH][BE], which is due to a deeper LUMO orbital over [BE] anion.

3.3. (IL)$_n$–B80 Systems: An Approach to the First Solvation Shell. Having discussed the features of the interaction mechanism between two [CH] based ILs and fullere B80, this section mainly pursues obtaining information
about the first IL solvation shell around of fullerene B$_{80}$ through the study of system composed by one fullerene and a variable number of ionic pairs ($n$). On the basis of the most stable configuration above-described, (IL)$_n$−B$_{80}$ systems ($n$ = 2, 3, 4 and 6) have been optimized. The binding energy as a function of $n$ was calculated as follows:

$$BE_{(IL)_n-B_{80}} = [(E_{B_{80}} + nE_{IL}) - E_{(IL)_n-B_{80}}]$$

(3)

where $E_{(IL)_n-B_{80}}$ is the total energy of (IL)$_n$−B$_{80}$ system. The evolution of $BE_{(IL)_n-B_{80}}$ is shown in Figure 5. According to this figure $BE_{(IL)_n-B_{80}}$ increases as a function of $n$ from 43.86 kcal mol$^{-1}$/45.95 kcal mol$^{-1}$ (for [CH]/[LAC]/[CH]/[BE]) at $n$ = 1 to 252.49 kcal mol$^{-1}$/262.16 kcal mol$^{-1}$ at $n$ = 6, with dispersion contributions lying between 16.22 kcal mol$^{-1}$/22.04 kcal mol$^{-1}$ and 119.39 kcal mol$^{-1}$/143.75 kcal mol$^{-1}$. Figure 5 also draws $BE_{(IL)_n-B_{80}}$ per ionic pair ($BE_{(IL)_n-B_{80}}/n$), whose average values per ionic pair are 43.18 and 44.13 kcal mol$^{-1}$ for [CH]/[LAC] and [CH]/[BE], respectively. These values per ionic pair are very close to that are computed for IL−B$_{80}$ systems, which suggest that IL−B$_{80}$ interactions is the main contribution to the total $BE_{(IL)_n-B_{80}}$. As expected, dispersion energy is one of the main contributions to the total energy, in average, dispersion contributions supply 43.34% and 51.45% to the total $BE_{(IL)_n-B_{80}}$ per ionic pair for [CH]/[LAC] and [CH]/[BE], respectively. For (IL)$_n$−B$_{80}$ systems, this slightly increase in the dispersion contribution could be to ion−ion interactions between different ionic pairs. Figures 6 and 7 illustrate optimized structure of (IL)$_n$−B$_{80}$ ($n$ = 2, 6) along RDG isosurfaces. As previously noted, their green colors also suggest that dispersion interactions play one of the main roles in the interaction between ILS and B$_{80}$. Figure 8 plots the evolution of the total charge over the cation, the anion and B$_{80}$ as a function of $n$. Although B$_{80}$ gains a negative charge equal to 0.07 e$^-$ for $n$ = 1, it charge becomes to be ≃ 0.0 from $n$ ≥ 3. In the same way, a positive/negative charge over the cation/anion also shows an asymptotic behavior with $n$, with a charge of
around +0.50 e−/−0.50 e−. Thus, B_{80} could be able to disturb charge distribution in the IL bulk.

4. CONCLUSIONS
In this work, DFT simulations were performed to obtain insights on the interaction mechanism and the approximation to the solvation of fullerene B_{80} with ILs. As a first approximation, solvation capability has been assessed through binding energies between selected ILs and B_{80} molecule for systems composed by one ionic pair and one fullerene B_{80}. Low toxic and biodegradable [CH] cation based ILs paired with [LAC] and [BE] anions were selected. DFT simulations results showed that a bond (which shows ionic character) is formed between COO− group in the anion with one boron atom in the fullerene B_{80}, which arises issues for the use of B_{80} dispersed in ILs. Nevertheless, our results also revealed useful information for the search of task-specific ILs for B_{80} solvation. It can be concluded that ions with aromatic motifs could be adequate as B_{80} solvents due to strong interactions through π-stacking. B_{80} atoms develop hydrogen bonding with nonaromatic [CH] cation, with boron atoms acting as hydrogen bond acceptors. Finally, information about the first solvation shell has been obtained through the optimization of (IL)−B_{80} systems with up to 6 ionic pair interacting with one fullerene B_{80}.

This work offers an insight up to the key parameters governing B_{80} solvation by ILs at the molecular level, providing some key ideas for the search of task specific solvent for B_{80}. Although the application of new system based on B_{80} ILs is still an unexplored field, a deep understanding of its behavior in solution (for instance dispersed in ILs) can increase its potential applications.

■ ASSOCIATED CONTENT

Supporting Information
The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcb.5b05187.

Detailed description of the conformational search procedure with examples (Figure S1) and density of states of B_{80} and IL−B_{80} systems (Figure S2) (PDF)

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Notes
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■ REFERENCES


Figure 8. Evolution of the total charge over the cation (red), the anion (green), and B_{80} fullerene (black). Total charges over ions are average value per ion. Points at n = 0 stands for isolated ionic liquids.


