# A SPECIAL SYSTEM OF BOUNDARY VALUE PROBLEMS 

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#### Abstract

In this paper we introduce a special system of boundary value problems and give a method for solving it. Then we give a detailed application of this method to a system of boundary value problems of Hilbert type.


Let $G_{1}$ and $G_{2}$ be simply connected bounded regions with simple closed contours $c_{1}$ and $c_{2}$ in the $z-$ plane and let $\bar{G}_{1} \cap \bar{G}_{2}=\Phi$.

The aim is to find two sectionally holomorphic functions $\varnothing_{1}(\mathrm{z})$ and $\varnothing_{2}(\mathrm{z})$ whose boundery values $\varnothing_{1}^{ \pm}(\mathrm{t})$ and $\varnothing_{2}^{ \pm}(\mathrm{t})$ satisfy the following conditions:

On $\mathrm{c}_{1}$

$$
\begin{equation*}
\mathrm{A}_{1}\left[\varnothing_{1}^{+}(\mathrm{t}), \varnothing_{1}^{-}(\mathrm{t})\right]=\mathrm{F}_{1}\left[\varnothing_{2}(\mathrm{t}), \overline{\varnothing_{2}(\mathrm{t})}\right] \tag{1}
\end{equation*}
$$

and on $\mathrm{c}_{2}$

$$
\begin{equation*}
A_{2}\left[\varnothing_{2}^{+}(t), \varnothing_{2}^{-}(t)\right]=F_{2}\left[\varnothing_{1}(t), \overline{\varnothing_{1}(t)}\right] \tag{2}
\end{equation*}
$$

where $A_{1}$ and $A_{2}$ are linear with respect to their arguments.
First we suppose that both of right hand sides of (1) and (2) are given, then from (1) and (2) we obtain $\varnothing_{1}$ and $\varnothing_{2}$ respectively.

## Boundary value problems

Consider that

$$
\begin{equation*}
\varnothing_{1}(z)=\frac{1}{2 \pi i} \int_{c_{1}} \frac{\varphi_{1}(\tau)}{\tau-z} . d t \tag{3}
\end{equation*}
$$

where the density function $\boldsymbol{\varphi}_{1}(\mathrm{t}) \in \mathrm{H}$ (Hölder class).
The boundary values of the function $\varnothing_{1}(z)$ take the form

$$
\left.\begin{array}{l}
\varnothing_{1}^{+}(t)=\frac{\varphi_{1}(t)}{2}+\frac{1}{2 \pi i} \int_{c_{1}} \frac{\varphi_{1}(t)}{\tau-t} d \tau \\
\varnothing_{1}^{-}(t)=-\frac{\varphi_{1}(t)}{2}+\frac{1}{2 \pi i} \int_{c 1} \frac{\varphi_{1}(t)}{\tau-t} d \tau \tag{4}
\end{array}\right\}
$$

Substituting from (4) into (1), we have

$$
\begin{align*}
& A_{1}\left[\frac{\varphi_{1}(t)}{2}+\frac{1}{2 \pi i} \int_{c_{1}} \frac{\varphi_{1}(\tau)}{2} d \tau,-\frac{\varphi_{1}(t)}{2}+\right. \\
& \left.\quad+\frac{1}{2 \pi i} \int_{c_{1}} \frac{\varphi_{1}(t)}{\tau-t} d \tau\right]=F\left[\varnothing_{2}(t), \overline{\varnothing_{2}(t)}\right] \ldots \ldots \tag{5}
\end{align*}
$$

and therefore we have the singular integral equation (5) with respect to the unknown function $\varphi_{1}(t)$ for which we can apply Noether's theorems. Under certain conditions we obtain $\varphi_{1}(t)$ and consequently $\varnothing_{1}(z)$.

It is known that the function $\varnothing_{1}(\mathrm{z})$ can be written in the form

$$
\begin{equation*}
\varnothing_{1}(z)=\frac{1}{2 \pi i} \int_{c_{1}} \frac{\epsilon_{1}\left[\varnothing_{2}(t), \overline{\varnothing_{2}(t)}\right]}{t-z} d t \tag{6}
\end{equation*}
$$

Thus, on $\mathrm{c}_{2}$

$$
\begin{equation*}
\varnothing_{1}(t)=\frac{1}{2 \pi i} \int_{c_{1}} \frac{\epsilon_{1}\left[\varnothing_{2}(\tau), \overline{\left.\varnothing_{2}(\tau)\right]}\right.}{\tau-t} d \tau \tag{7}
\end{equation*}
$$

Subtituting from (7) into (2), we have

$$
\begin{align*}
& A_{2}\left[\varnothing_{2}^{+}(t), \varnothing_{2}^{-}(t)\right]=F_{2}\left[\frac{1}{2 \pi i} \int_{c_{1}} \frac{\epsilon_{1}\left[\varnothing_{2}(\tau), \overline{\varnothing_{2}(\tau)}\right]}{\tau-t} d \tau\right. \\
& \left.\quad-\frac{1}{2 \pi i} \int_{c_{1}} \frac{\overline{\epsilon_{1}\left[\varnothing_{2}(t), \overline{\varnothing_{2}(t)}\right]}}{\bar{\tau}-\bar{t}} \overline{d \tau}\right] \ldots \ldots \ldots \ldots \ldots \ldots \tag{8}
\end{align*}
$$

Let $A_{2}$ have properties such that $\varnothing_{2}(z)$ can be obtained from (8) and assume that whenever $z=t$, on $c_{1}$, we define $\varnothing_{2}(t)$.

Thereby from (8), we immediately obtain an integral equation with respect to $\varnothing_{2}(\mathrm{t})$.
By obtaining the solution $\varnothing_{2}(t)$ of such integral equation the function $\varnothing_{1}(\mathrm{z})$
follows directly from (6). Similarly, we find $\varnothing_{2}(z)$.
The method is complete.

We now give an application of this method:
Consider the following system of boundary value problems of Hilbert type [1].

On $c_{1}$

$$
\begin{equation*}
\varnothing_{1}^{+}(t)-A_{1}(t) \varnothing_{1}^{-}(t)=f_{1}(t)+\propto_{1}(t) \varnothing_{2}(t)+\propto_{2}(t) \overline{\varnothing_{2}(t)} \tag{9}
\end{equation*}
$$

and on $\mathrm{c}_{2}$

$$
\begin{equation*}
\varnothing_{2}^{t}(t)-A_{2}(t) \varnothing_{2}^{-}(t)=f_{2}(t)+\beta_{1}(t) \varnothing_{1}(t)+\beta_{2}(t) \overline{\varnothing_{1}(t)} \tag{10}
\end{equation*}
$$

and let the index ae of $A_{1}(t)$ be not negative.

From (9), we have

$$
\begin{align*}
\varnothing_{1}(z)=\frac{X(z)}{2 \pi i} \int_{c_{1}} & \frac{f_{1}(t)+\propto_{1}(t) \varnothing_{2}(t)+\propto_{2}(t) \overline{\varnothing_{2}(t)}}{X^{+}(t)} \frac{d t}{t-z}+ \\
& +p_{a_{1}}(z) X(z) \ldots \ldots \ldots \ldots \ldots \ldots \tag{11}
\end{align*}
$$

where $X(z)$ is the canonical function of the associated homogeneous equation with respect to (9) and $p_{\mathrm{ae}_{1}}(z)$ is a polynomial of degree $a e_{1}$ with arbitrary coefficients.

Whenever $z=t$, on $c_{2}$, then

$$
\begin{align*}
\varnothing_{1}(t)=\frac{X(t)}{2 \pi i} \int_{c_{1}} & \frac{f_{1}(\tau)+\propto_{1}(\tau) \varnothing_{2}(\tau)+\propto_{2}(\tau) \overline{\varnothing_{2}(\tau)}}{X^{+}(\tau)} \frac{d \tau}{\tau-t} \\
& +p_{\mathrm{ae}_{1}}(t) X(t) \cdots \cdots \cdots \cdots \cdots \cdots \cdots \tag{12}
\end{align*}
$$

Thus, the right hand side of $(10)$ can be written in the form

$$
\begin{array}{r}
F(t), \frac{1}{2 \pi i} \cdot \int_{c_{1}} \frac{\propto_{1}(\tau) \varnothing_{2}(\tau)+\propto_{2}(\tau) \overline{\varnothing_{2}(\bar{\tau})}}{X^{+}(\tau)} \frac{d \tau}{\tau-t} \\
\frac{1}{2 \pi i} \int_{c_{1}} \frac{\overline{\alpha_{1}(\tau)} \overline{\varnothing_{2}(\tau)}+\overline{\propto_{2}(\tau)} \varnothing_{2}(\tau)}{X^{+}(\tau)} \frac{d \bar{\tau}}{\overline{\tau-\bar{t}}} \\
=F \varnothing_{2} \ldots \ldots \ldots \ldots \ldots \ldots \ldots \tag{13}
\end{array}
$$

From (13) the boundary condition (10)

$$
\begin{equation*}
\varnothing_{2}^{+}(t)-A_{2}(t) \varnothing_{2}^{-}(t)=F \varnothing_{2} \tag{14}
\end{equation*}
$$

Let the index $\mathrm{ae}_{2}$ of $\mathrm{A}_{2}(t)$ be not negative. Then we have

$$
\begin{equation*}
\varnothing_{2}(t)=\frac{y(t)}{2 \pi i} \int_{C_{2}} \frac{F \varnothing_{2}}{y^{+}(t)} \frac{d \tau}{\tau-z}+Q_{a_{2}}(z) y(z) . \tag{15}
\end{equation*}
$$

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where $y(z)$ is the cononical function of the associated homogeneous equation with respect to (14) and $Q_{a_{2}}$ with arbitrary coefficients.

Whenever $z=t$, on $c_{1}$, we have

$$
\begin{equation*}
\varnothing_{2}(t)=\frac{y(t)}{2 \pi i} \int_{c_{2}} \frac{F \varnothing_{2}}{y^{+}(\tau)} \frac{d \tau}{\tau-t}+Q_{\mathrm{ae}_{2}}(t) y(t) \tag{16}
\end{equation*}
$$

and therefore from (13) and (16), we immediately obtain an integral equation with respect to $\varnothing_{2}(t)$. By applying Fredholm's Integral Equation Theory we obtain $\varnothing_{2}(t)$ and consequently from (11) we find $\varnothing_{1}(z)$. Similarly we find $\varnothing_{2}(\mathrm{z})$.

## REFERENCES

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## نظام خاص من مسائل القيير الحدية <br> محمد إبراهيم محمد

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