

A SHORT NOTE ON THE POINT-WISE SUMMABILITY OF THE CONJUGATE SERIES OF A FOURIER SERIES IN THE NORLUND SENSE

By
ZIAD RUSHDI ALI
United Arab Emirates

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ABSTRACT

In this note we consider an analogous theorem of [1] in the sense of point-wise summability of the conjugate series of a Fourier series in the Norlund sense

INTRODUCTION

1. Let $f(t)$ be a periodic function with period 2π , integrable in the sense of Lebesgue over $(-\pi, \pi)$. Let $\tilde{S} [f]$ denote the conjugate series of the Fourier series of f . Let $\phi (t) = f(x+t) + f(x-t) - 2 f(x)$, $\psi (t) = f(x+t) - f(x-t)$,

$$\Psi (t) = \int_0^t |\psi (u)| du, \text{ and } \Phi (t) = \int_0^t |\phi (u)| du$$

In[1] we have :

Theorem 1.1 : Let $q_n = \frac{1}{(n+1)^\alpha}$, $0 \leq \alpha < 1$. Then $S [f]$ is summable $S(N, q_n)$ to $f(x)$ at each point x where $\Phi (t) = o(t)$.

(The case $\alpha = 0$ is Lebesgue's theorem [5])

In this note we prove the following analogous theorem of theorem 1.1 above.

Theorem 1: Let $q_n = \frac{1}{(n+1)^\alpha}$, $0 \leq \alpha < 1$. Then $\tilde{S} [f]$ is

summable $S(N, q_n)$ to $\frac{1}{\pi} \int_0^\pi \frac{\psi (t)}{2 \tan \frac{1}{2} t}$ at each point x

where $\Psi (t) = o(t)$.

Proof: Let $\tilde{S}_n (x)$ denote the sequence of partial sums of the conjugate series of a Fourier series, and let $t_n(x)$ be the corresponding Norlund mean. Then clearly ([3] , and [5])

$$t_n(x) = \sum_{k=0}^n q_k \frac{1}{2\pi Q_n} \int_0^\pi \psi(t) \frac{\cos \frac{t}{2} \cos (n-k + \frac{1}{2}) t}{\sin \frac{t}{2}} dt$$

Hence

$$t_n(x) - \frac{1}{\pi} \int_0^\pi \frac{\psi(t)}{2 \tan \frac{t}{2}} dt = - \int_0^\pi \psi(t) K_n(t) dt, \text{ where}$$

$$K_n(t) = \frac{1}{2\pi Q_n} \sum_{k=0}^n \frac{\cos (n-k + \frac{1}{2}) t}{\sin \frac{t}{2}}.$$

In order to prove the theorem we show that :

$$\int_0^\pi \psi(t) K_n(t) dt = o(1) \quad \text{as} \quad n \rightarrow \infty.$$

Now

$$\begin{aligned} \int_0^\pi \psi(t) K_n(t) dt &= \left(\int_0^{\frac{1}{n}} + \int_{\frac{1}{n}}^1 + \int_1^{\widehat{\pi}} \right) \psi(t) K_n(t) dt \\ &= \widetilde{I}_1 + \widetilde{I}_2 + \widetilde{I}_3, \text{ say} \end{aligned}$$

First by [2] above we have:

$$\begin{aligned} \widetilde{I}_1 &= \int_0^{\frac{1}{n}} \psi(t) K_n(t) dt \\ &= o(n \int_0^{\frac{1}{n}} |\psi(t)| dt) \\ &= o(1) \text{ as } n \rightarrow \infty. \end{aligned}$$

Second, clearly the method $S(N, \frac{1}{(n+1)^\alpha})$ is regular. Hence by

the Riemann Lebesgue theorem, and the regularity of the method $S(N, \frac{1}{(N+1)^\infty})$ we have :

$$\widetilde{I}_3 = \int_\pi^1 \psi(t) K_n(t) dt = o(1) \quad \text{as } n \rightarrow \infty.$$

Third by [4] above we have:

$$\bar{I}_2 = \int_{\frac{1}{n}}^1 \psi(t) K_n(t) dt = o\left(\frac{1}{Q_n} \int_{\frac{1}{n}}^1 |\psi(t)| \frac{Q_\tau}{t} dt\right).$$

Now

$$\frac{1}{Q_n} \int_{\frac{1}{n}}^1 |\psi(t)| \frac{Q_\tau}{t} dt = \left(\int_{\frac{1}{n}}^1 + \int_{\frac{1}{n-1}}^{\frac{1}{n-2}} + \dots + \int_{\frac{1}{2}}^1\right) |\psi(t)| \frac{Q_\tau}{t} dt$$

Hence by integration by parts and simplifying we obtain:

$$\frac{1}{Q_n} \int_{\frac{1}{n}}^1 \psi(t) \frac{Q_\tau}{t} dt - o(1) = \frac{1}{Q_n} \left\{ \Psi(t) \frac{Q_\tau}{t} \Big|_{\frac{1}{n}}^1 + \psi(t) \frac{Q_\tau}{t^2} dt \right\}$$

Now

$$\frac{1}{Q_n} \Psi(t) \frac{Q_\tau}{t} \Big|_{\frac{1}{n}}^1 = O\left(\frac{1}{Q_n}\right) + o\left(\frac{1 \cdot 1 \cdot Q_n}{Q_n \cdot n \cdot 1}\right) = o(1) \text{ as } n \rightarrow \infty$$

$$\begin{aligned} \frac{1}{Q_n} \int_{\frac{1}{n}}^1 \psi(t) \frac{Q_\tau}{t^2} dt &= \frac{1}{Q_n} \int_{\frac{1}{n}}^1 \Psi\left(\frac{1}{u}\right) Q(u) du \\ &= o\left(\frac{1}{Q_n} \int_{\frac{1}{n}}^1 \frac{Q(u)}{u} du\right) \end{aligned}$$

$$= o(1) \text{ as } n \rightarrow \infty \text{ by [1]}$$

This completes the proof of theorem 1.

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